

Large-eddy simulation of very stable boundary layers

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August, 2018

Motivation

Simulate **very** stable conditions using the buoyancy-adjusted stretch-vortex SGS model in LES with a novel LES model setup:

Surface fluxes dynamically computed by partially resolving the surface layer

Study of the:

1. SBL structure in very stable conditions with weak geostrophic winds
2. Turbulence structure and its energetic regimes

Large-eddy simulations

Initial and boundary conditions of second set of simulations

- $u_g = 1$ and 2 m s^{-1} (meridional component set to zero).
- Initial wind profile: Ekman solution.
- Initial potential temperature profile: uniform and equal to the reference temperature of 275 K up to 2 m , with an overlying stable lapse rate of 0.01 K m^{-1} .
- Surface cooling rates equal to 1 and 3 K h^{-1} .
- Surface fluxes dynamically computed by partially resolving the surface layer.
- Flat and smooth surface.
- Latitude: 70°N .
- Uniform and isotropic grid.
- Rayleigh damping layer imposed at the top of the domain (20%).
- Runs were carried out for 24 h of physical time (i.e., no diurnal cycle).

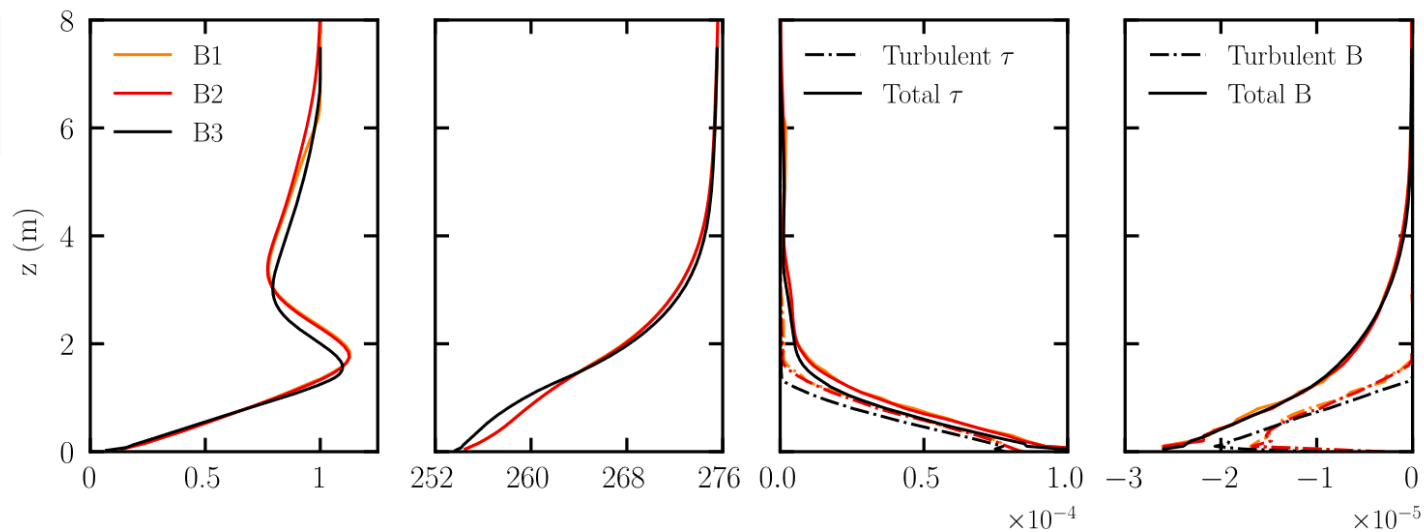
Large-eddy simulations

Table 2.2 Summary of the second set of simulations, where the surface fluxes were dynamically computed. “Run” shows the name of the simulation followed by geostrophic wind (u_g), surface cooling rate (C_r), number of grid points, grid resolution that is isotropic (Δ), surface friction velocity (u_*), surface heat flux (Q_*), Obukhov length (L), and height of the SBL (h). The bulk quantities correspond to a time average in $t = 14 - 15$ h.

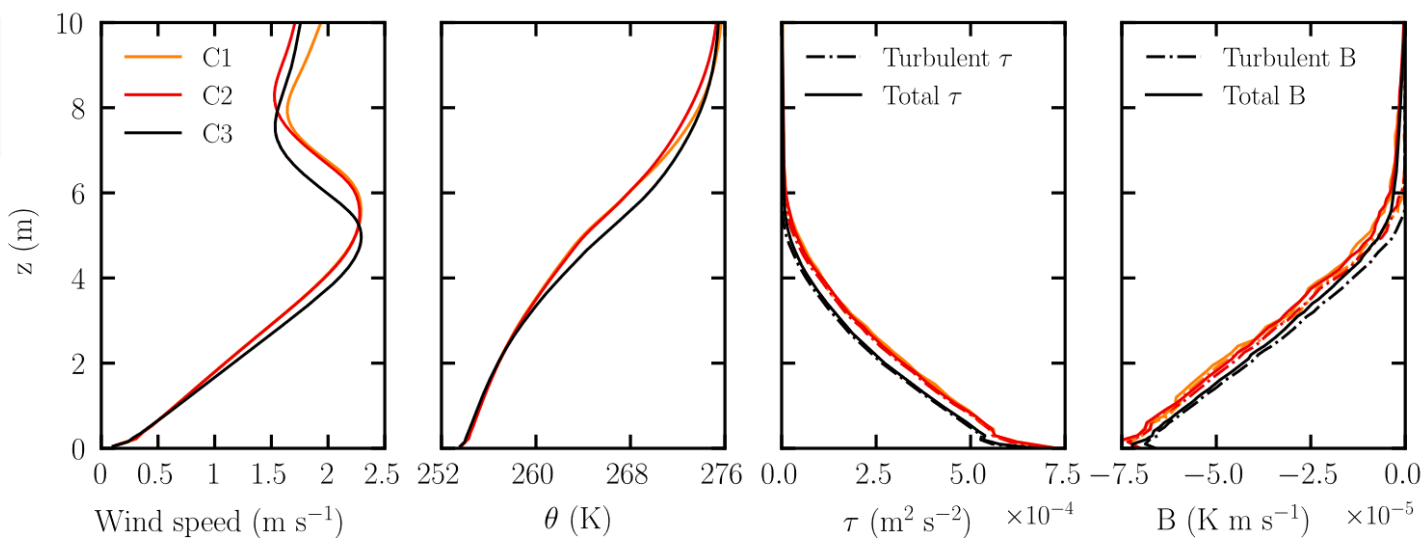
Run	u_g (m s^{-1})	C_r (K h^{-1})	Grid points	Δ (m)	$u_* \times 10^{-2}$ (m s^{-1})	$Q_* \times 10^{-3}$ (K m s^{-1})	L (m)	h (h)
A1	1	1	$64^2 \times 80$	0.10	1.60	-0.519	0.55	4.00
A2	1	1	$128^2 \times 100$	0.10	1.59	-0.520	0.55	4.00
A3	1	1	$256^2 \times 150$	0.05	1.46	-0.453	0.48	3.20
B1	1	3	$64^2 \times 80$	0.10	1.11	-0.718	0.13	1.70
B2	1	3	$128^2 \times 100$	0.10	1.10	-0.716	0.13	1.70
B3	1	3	$256^2 \times 150$	0.05	1.08	-0.722	0.12	1.35
C1	2	3	$64^2 \times 100$	0.15	2.70	-2.20	0.64	6.45
C2	2	3	$128^2 \times 120$	0.15	2.70	-2.10	0.64	6.30
C3	2	3	$256^2 \times 200$	0.10	2.59	-2.00	0.60	5.50

Sensitivity of the LES results

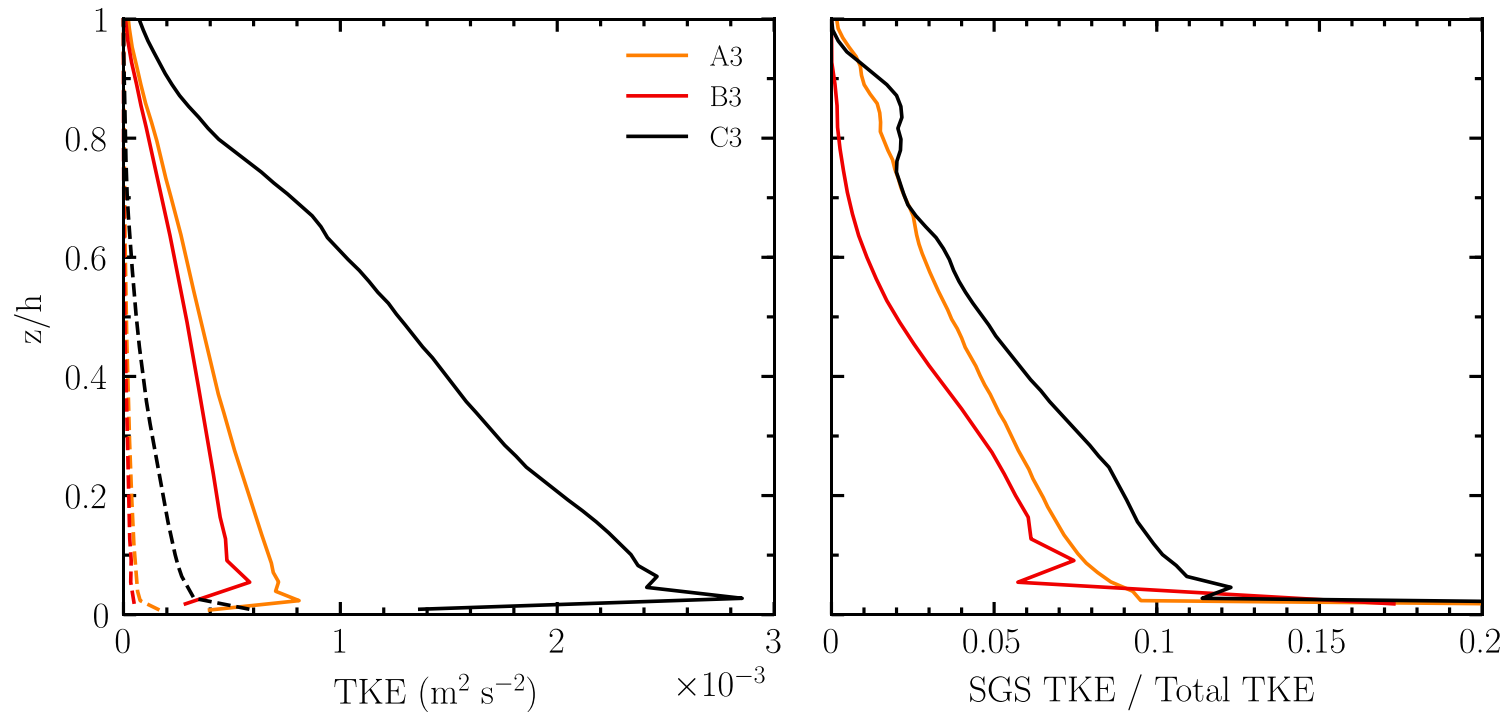
$\Delta_{B1} = \Delta_{B2} = 0.10$ m
 $\Delta_{B3} = 0.05$ m



$\Delta_{C1} = \Delta_{C2} = 0.15$ m
 $\Delta_{C3} = 0.10$ m



Sensitivity of the LES results – Grid convergence



Overall, more than **90%** of TKE is **resolved**.

A: u100.100
B: u100.300
C: u200.300

SBL height estimation methods

Table 2.3 Boundary layer height estimated by the following five methods: (1) height at which the vertical turbulent momentum flux vanishes; (2) height of the low-level jet; (3) height at which $\partial\bar{\theta}/\partial z$ is a maximum; and, (4) height at which the total TKE reduces to 5% of its surface value. The values correspond to a time average in $t = 14 - 15$ h.

Run	h_τ (m)	h_j (m)	$h_{\nabla\theta}$ (m)	h_E (m)
A	2.84	3.10	3.25	3.20
B	1.21	1.55	1.35	1.35
C	4.74	5.00	5.00	5.50
GABLS1	179	178	212	215

GABLS4_{0.25m} 20 21.5 17 22

$$h_E/L = 22/3.12 = 7$$

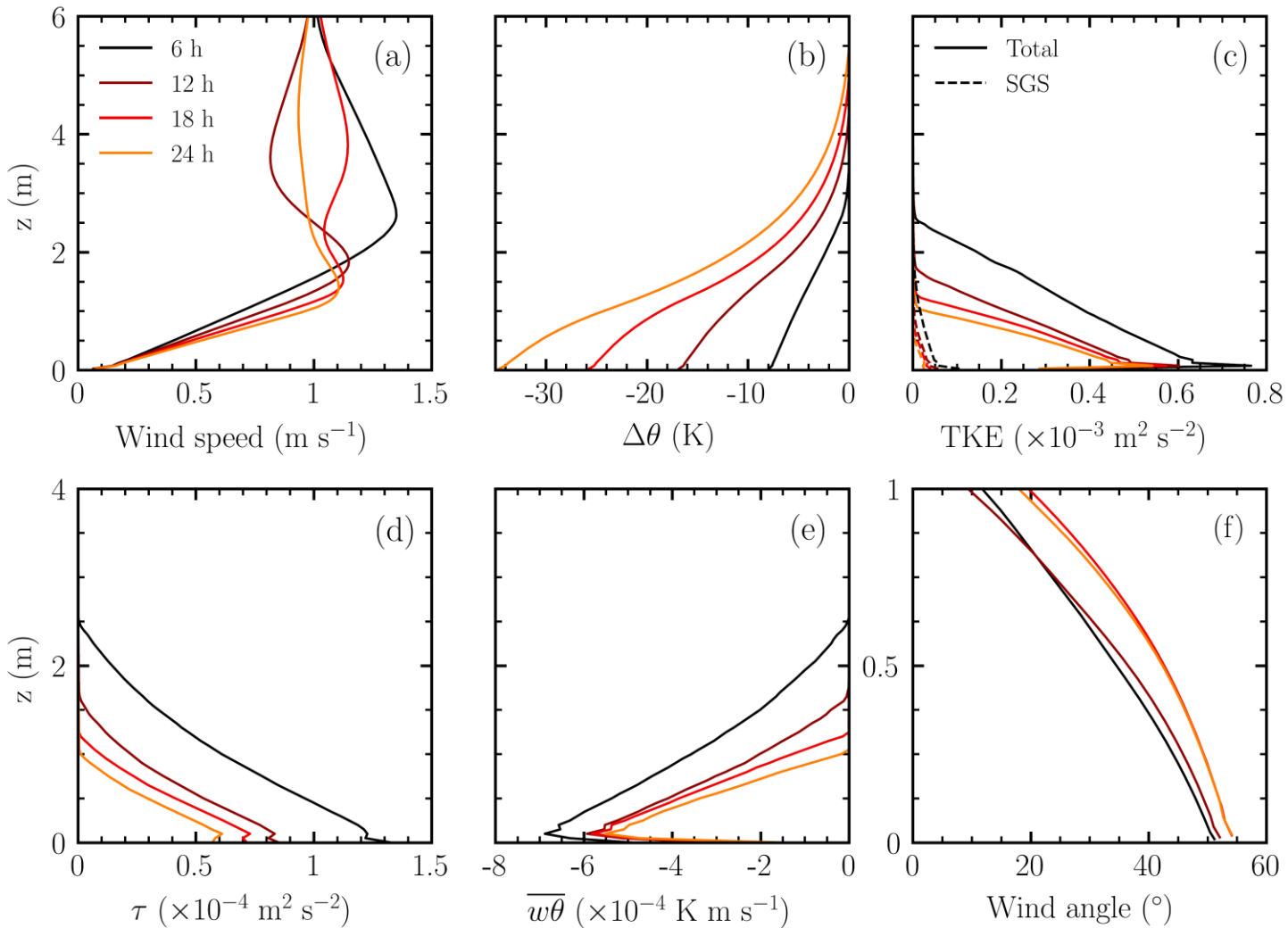
for $t = 15 - 16$ h

- h_E is chosen for this study.

A: u100.100
B: u100.300
C: u200.300

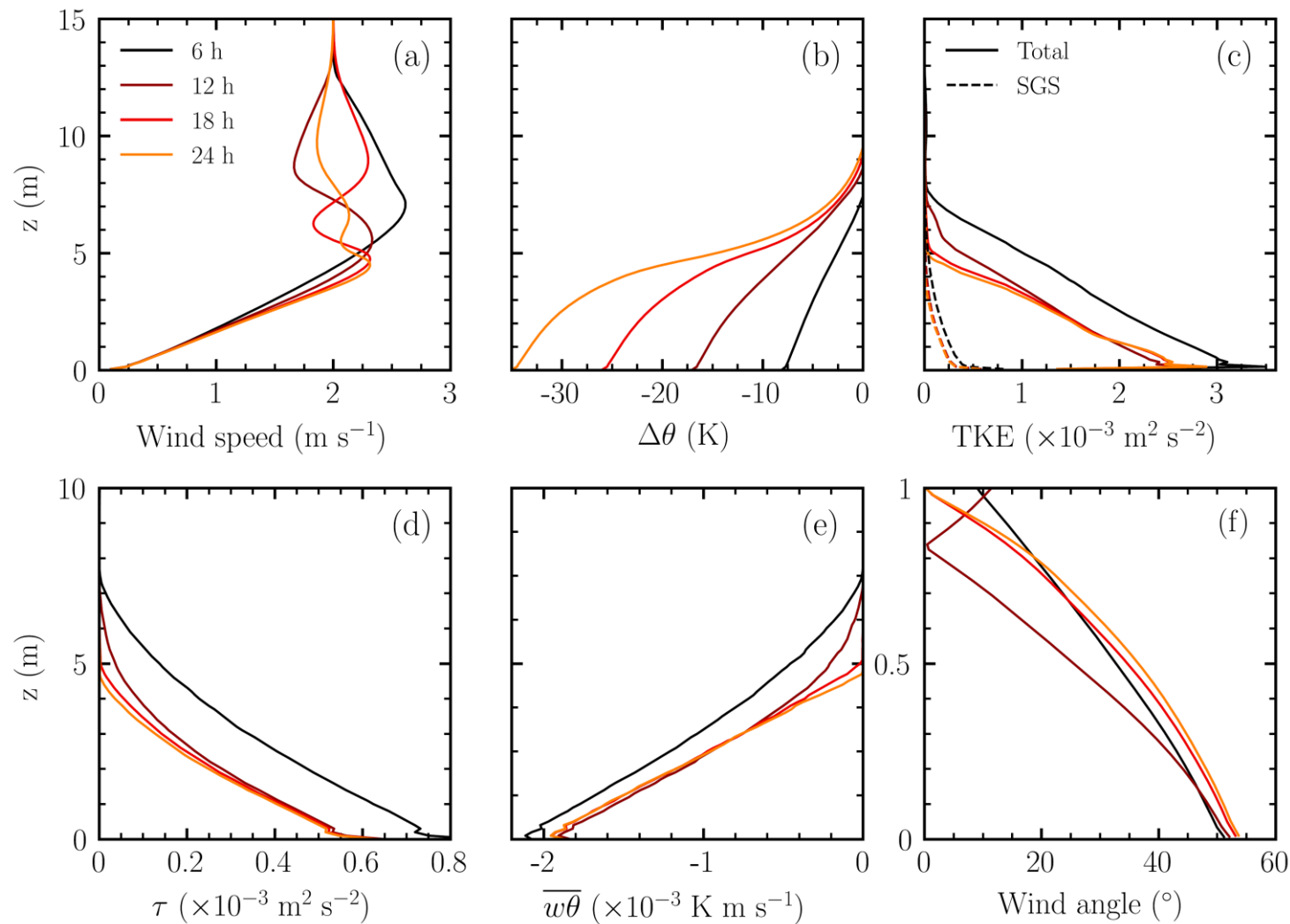
Temporal evolution of vertical profiles

Temporal evolution of simulation B ($u_g = 1 \text{ m s}^{-1}$ and $C_r = 3 \text{ K h}^{-1}$).



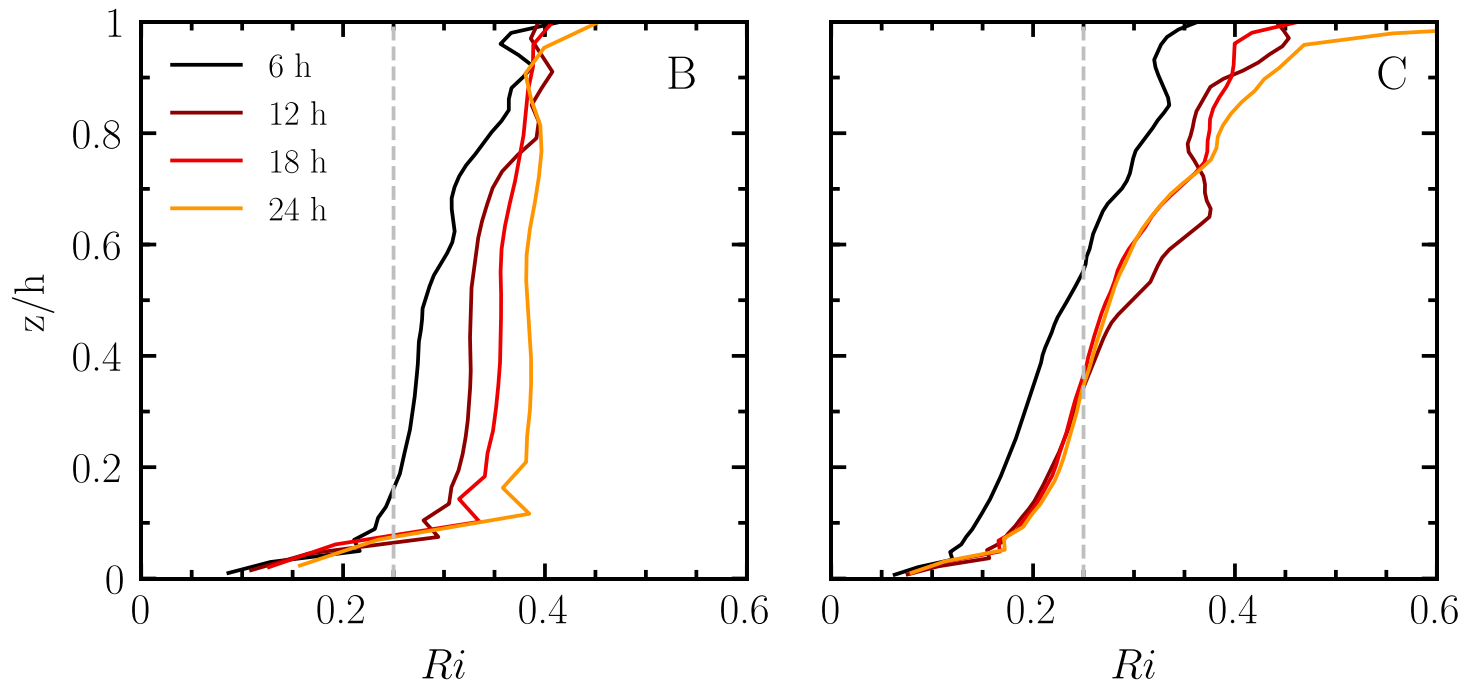
Temporal evolution of vertical profiles

Temporal evolution of simulation C ($u_g = 2 \text{ m s}^{-1}$ and $C_r = 3 \text{ K h}^{-1}$).

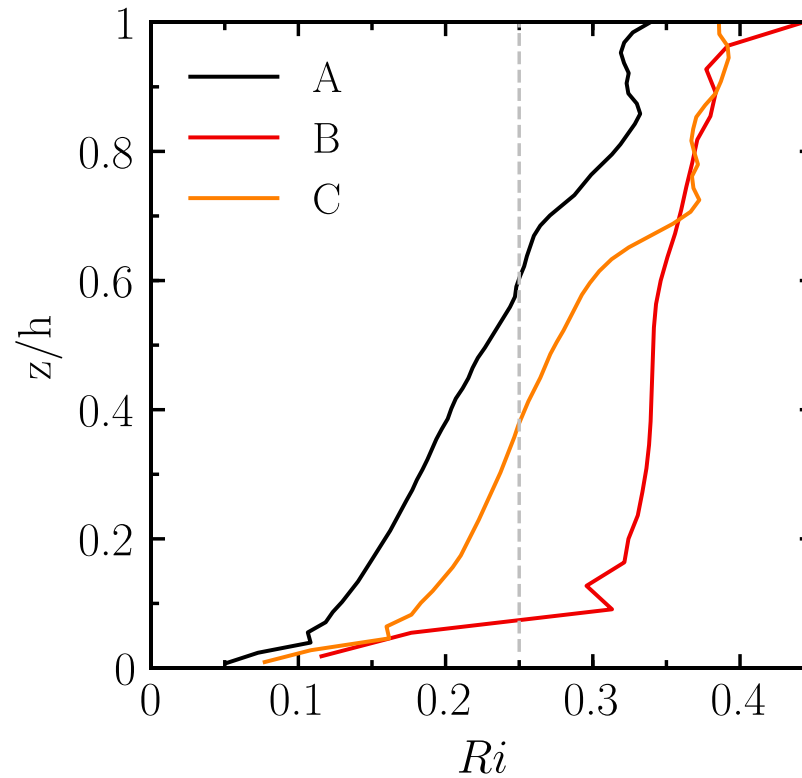


Temporal evolution of vertical profiles

Temporal evolution of Ri for simulations B ($u_g = 1 \text{ m s}^{-1}$ and $C_r = 3 \text{ K h}^{-1}$), and C ($u_g = 2 \text{ m s}^{-1}$ and $C_r = 3 \text{ K h}^{-1}$).



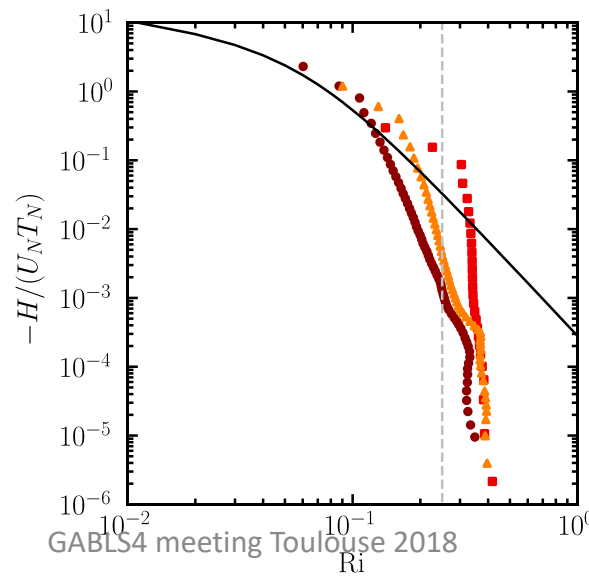
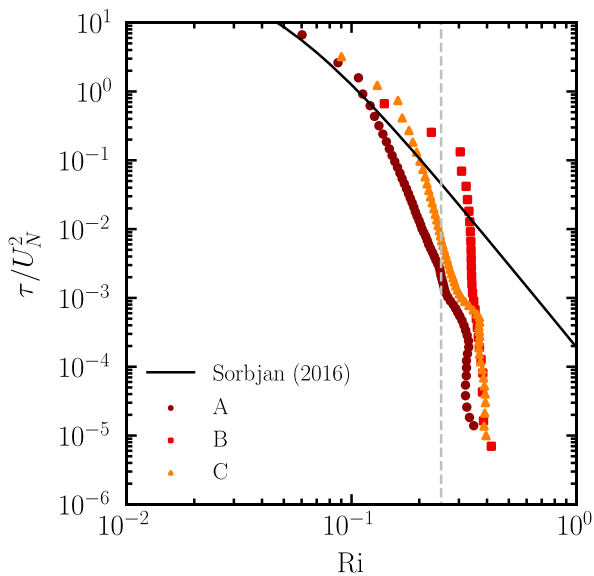
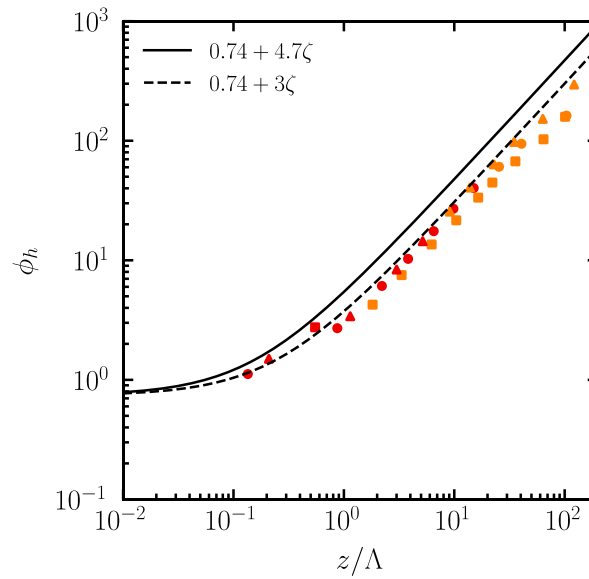
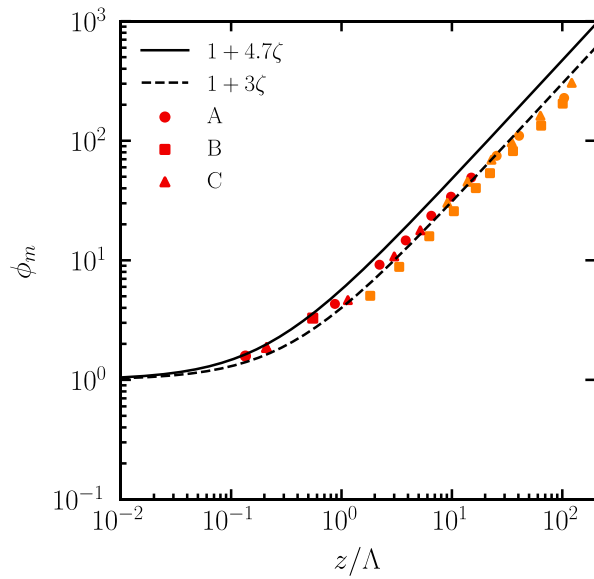
Gradient Richardson number



A: u100.100
B: u100.300
C: u200.300

Stability: $B > C > A$

Scaling systems



$$\frac{\tau}{U_N^2} = \frac{1}{Ri(1 + 300Ri^2)^{3/2}},$$

$$U_N = lN$$

$$\frac{-\overline{w'\theta'}}{U_N T_N} = \frac{1}{0.9Ri^{1/2}(1 + 250Ri^2)^{3/2}},$$

$$T_N = l \frac{\partial \bar{\theta}}{\partial z}$$

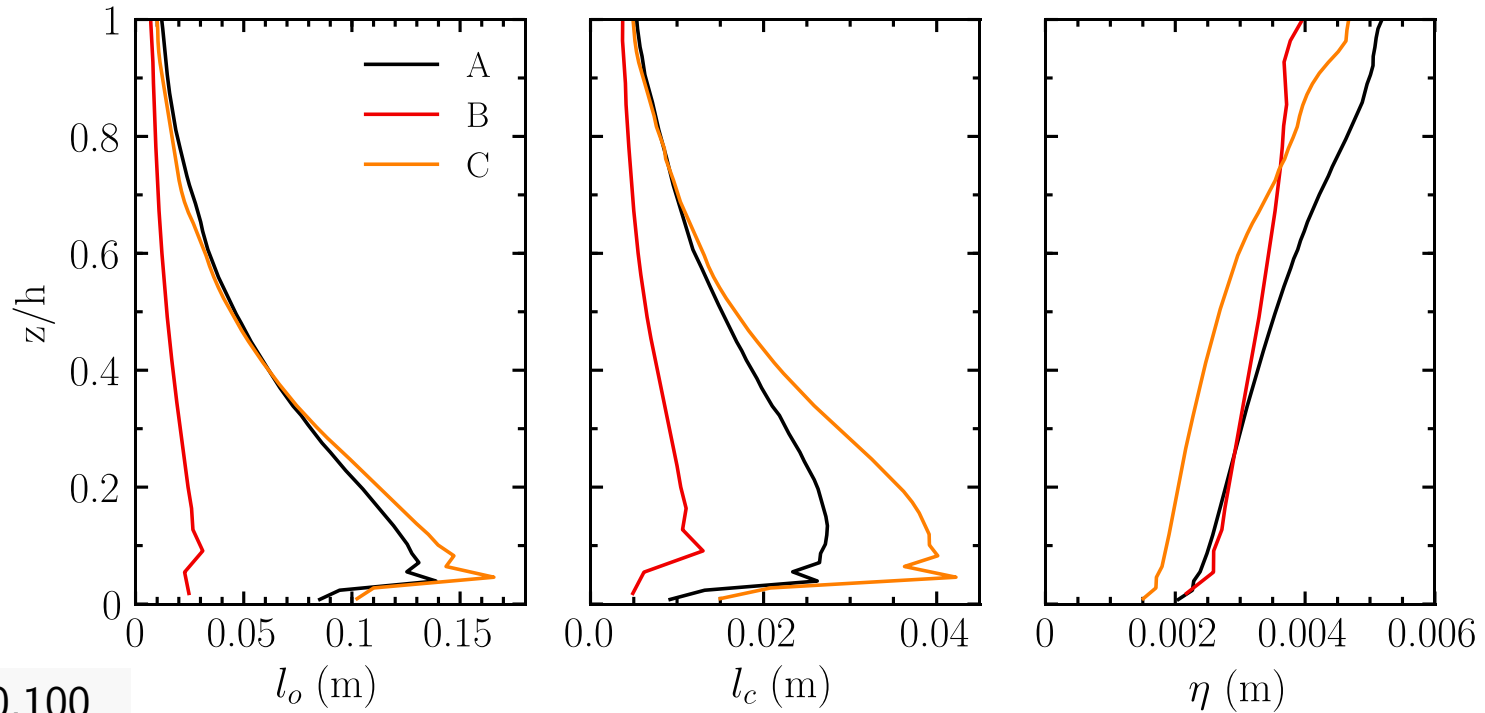
Length scales

- Dougherty-Ozmidov (l_o), Corrsin (l_c), and Kolmogorov (η) scales:

$$l_o = (\varepsilon/N^3)^{1/2}$$

$$l_c = (\varepsilon/S^3)^{1/2}$$

$$\eta = (v^3/\varepsilon)^{1/4}$$



A: u100.100

B: u100.300

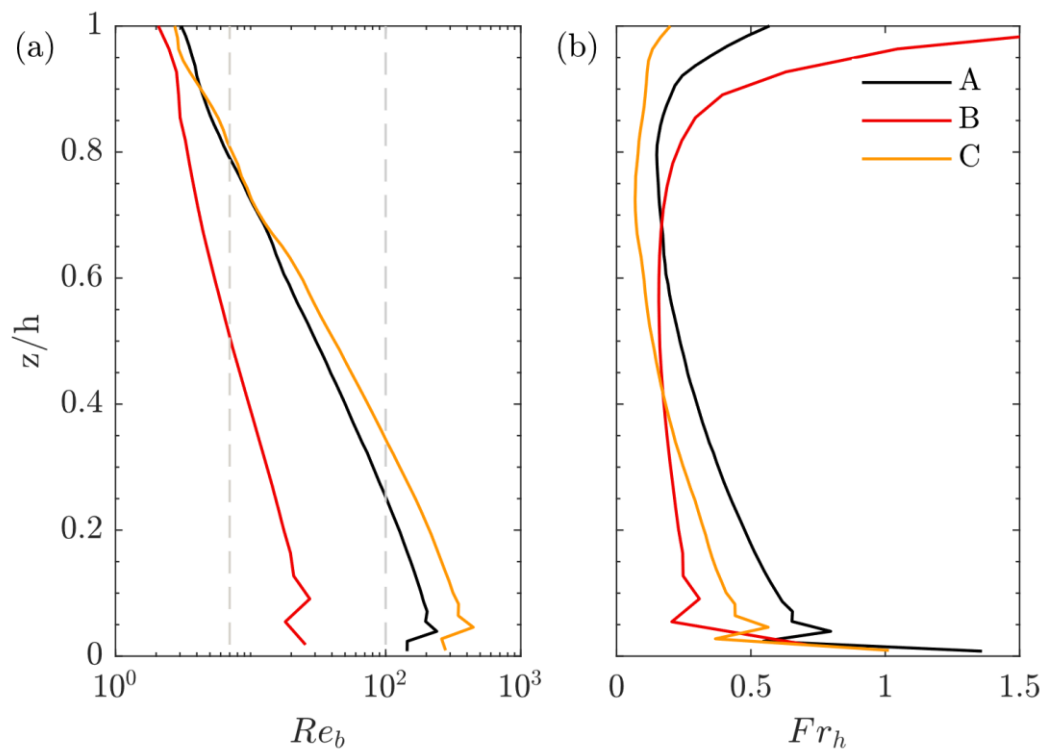
C: u200.300

Strongly stratified turbulence theory

- Strongly stratified turbulence regime without viscous effects (e.g. Billant and Chomaz 2001): $Re_b \gg 1$ and $Fr_h \ll 1$.

$$Re_b = \varepsilon / (\nu N^2) = (l_o / \eta)^{4/3}$$

$$Fr_h = \varepsilon / (Nu_{rms}^2) = (l_o / L_h)^{2/3}$$



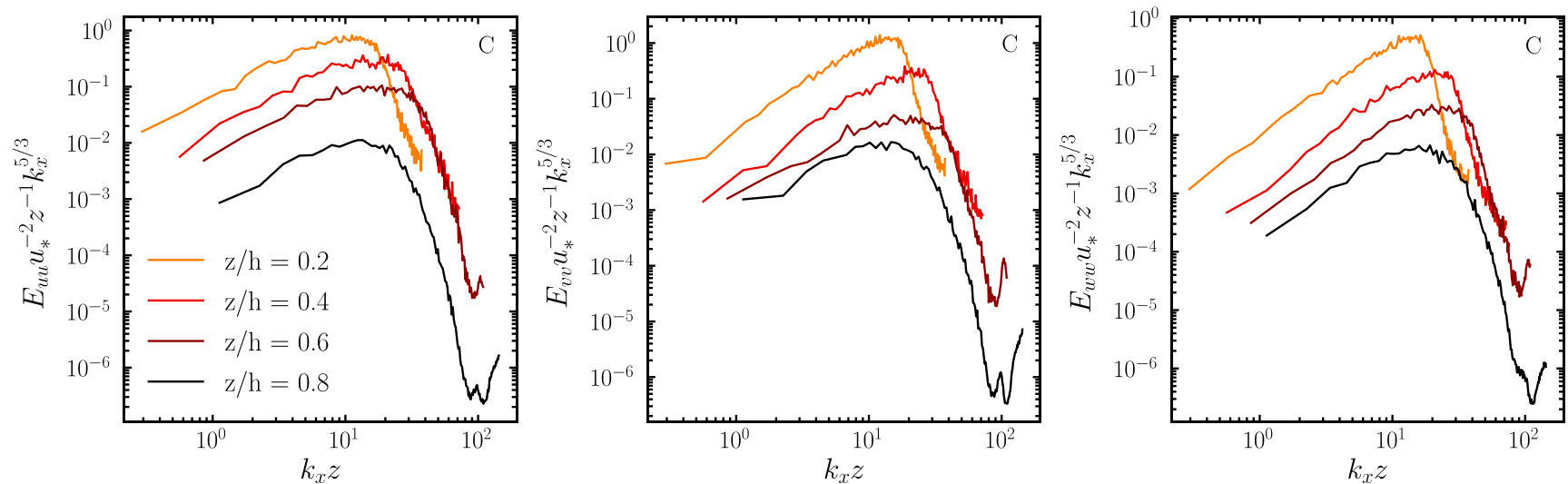
A: $u_{100.100}$

B: $u_{100.300}$

C: $u_{200.300}$

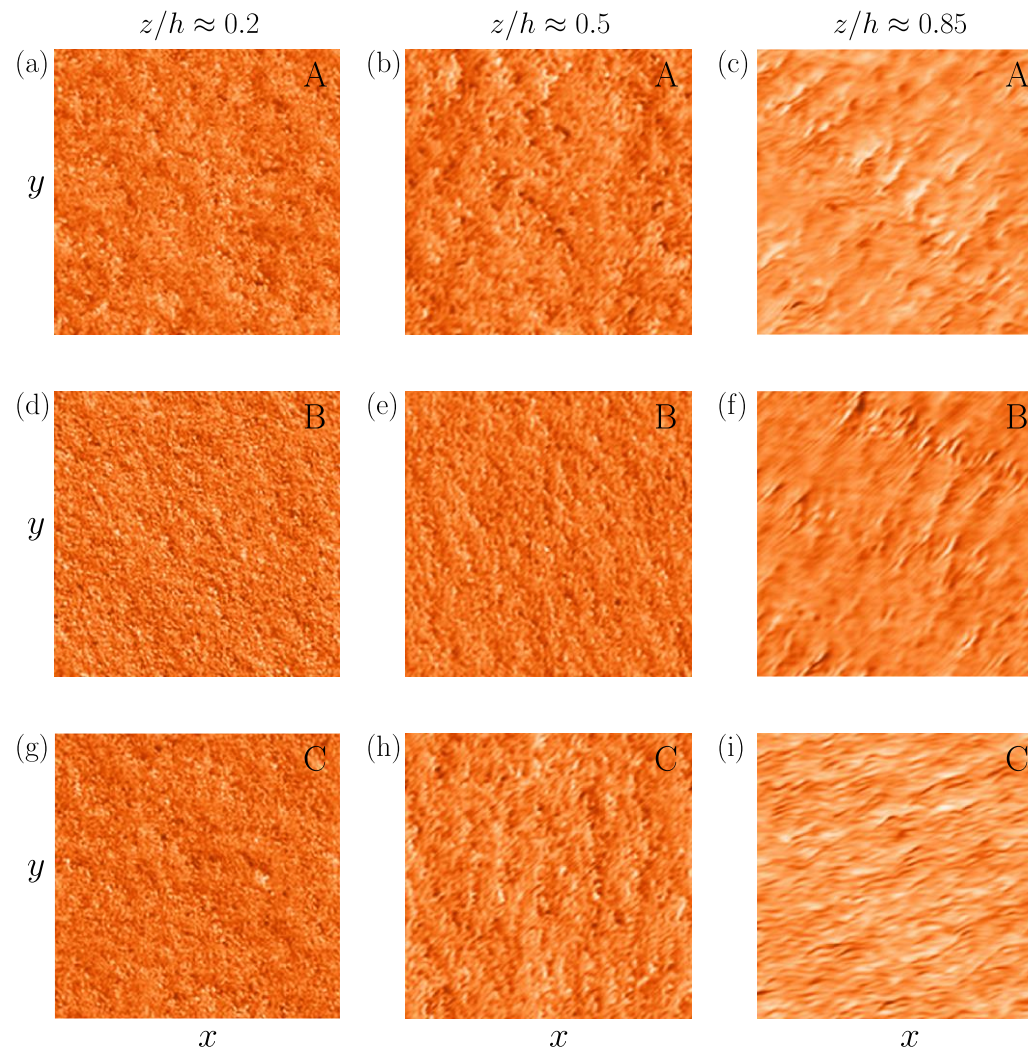
Spectra

- Compensated one-dimensional spectra of zonal, meridional and vertical wind components for simulation C (instant $t = 15$ h).



- $l_o < \Delta_x \rightarrow$ production, inertial and dissipation ranges correspond to anisotropic turbulence due to buoyancy effects.

Flow visualization



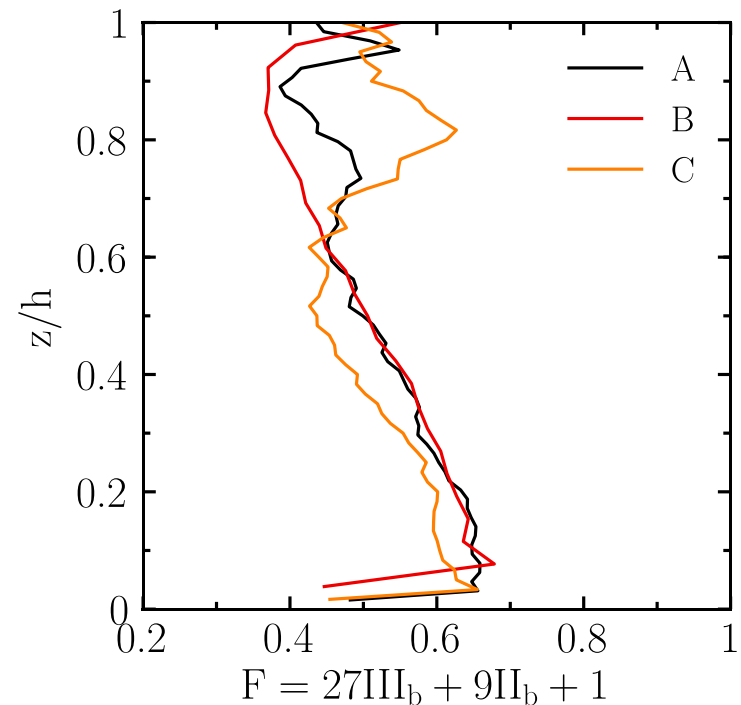
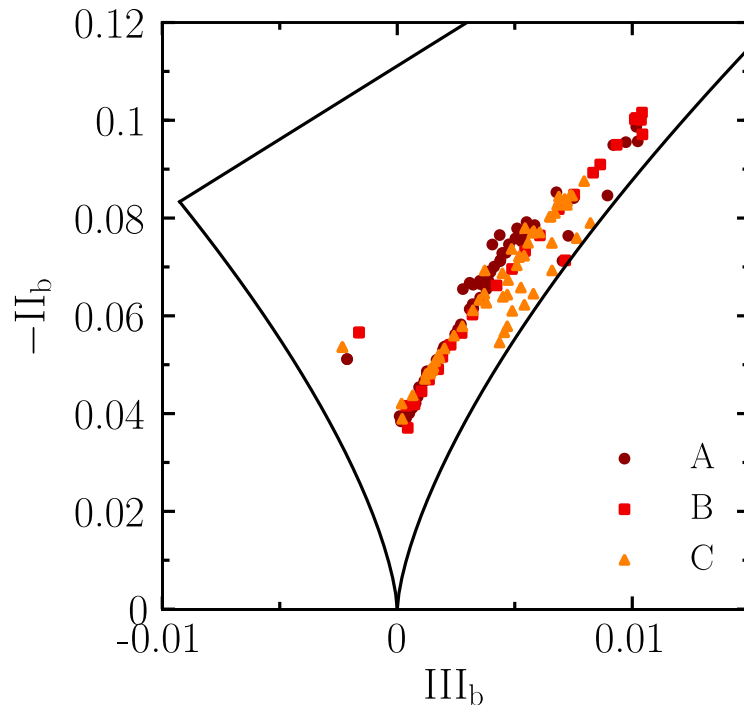
A: $u100.100$
B: $u100.300$
C: $u200.300$

Anisotropy invariant map

$$b_{ij} = \frac{\overline{u'_i u'_j}}{\overline{u'_k u'_k}} - \frac{\delta_{ij}}{3}$$

$$II_b = -\frac{1}{2} b_{ij} b_{ji}$$

$$III_b = \frac{1}{3} b_{ij} b_{jk} b_{ki}$$

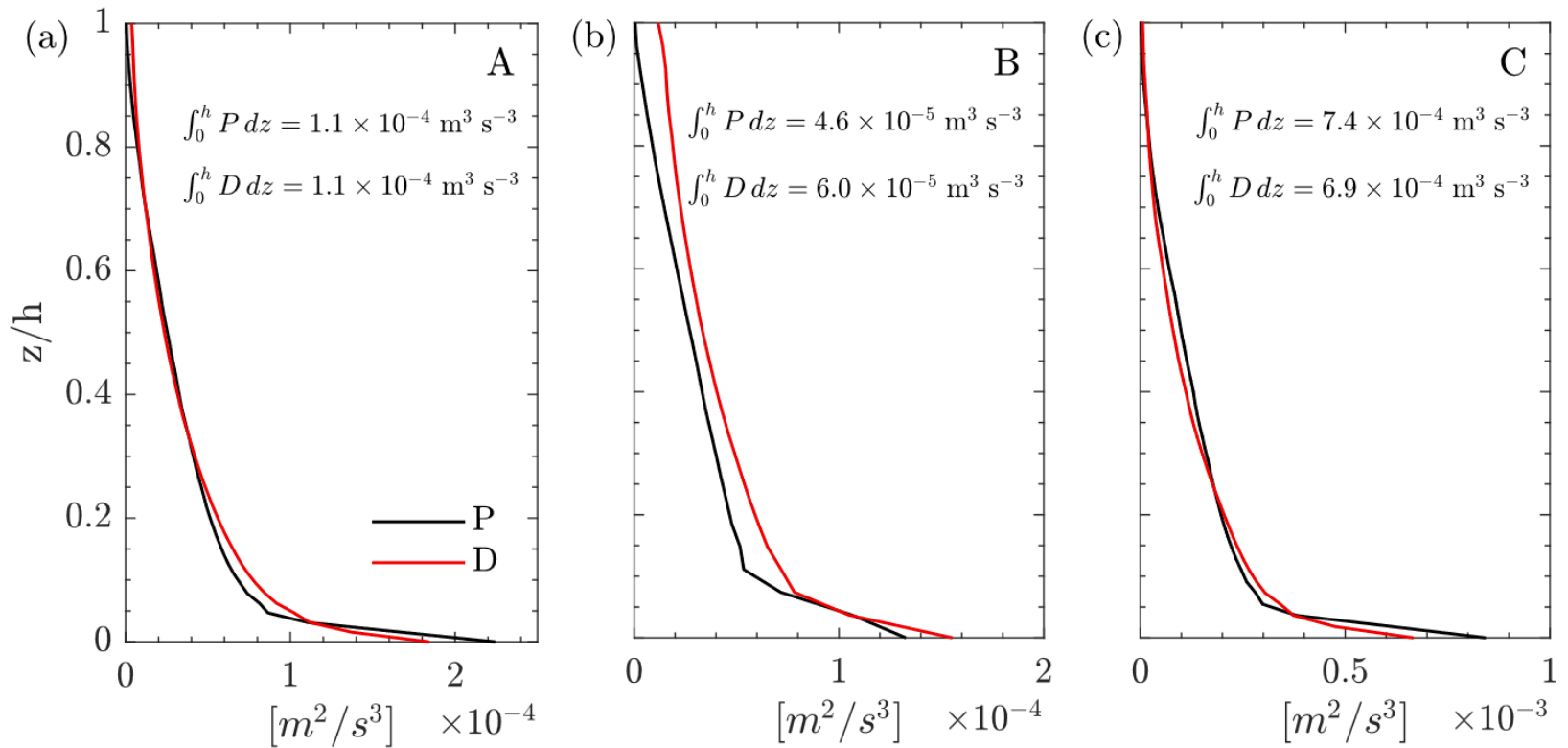


- Flow is dominated by one component $\overline{u'u'}$ (cigar-shaped turbulence).

TKE budget

P – shear production of TKE

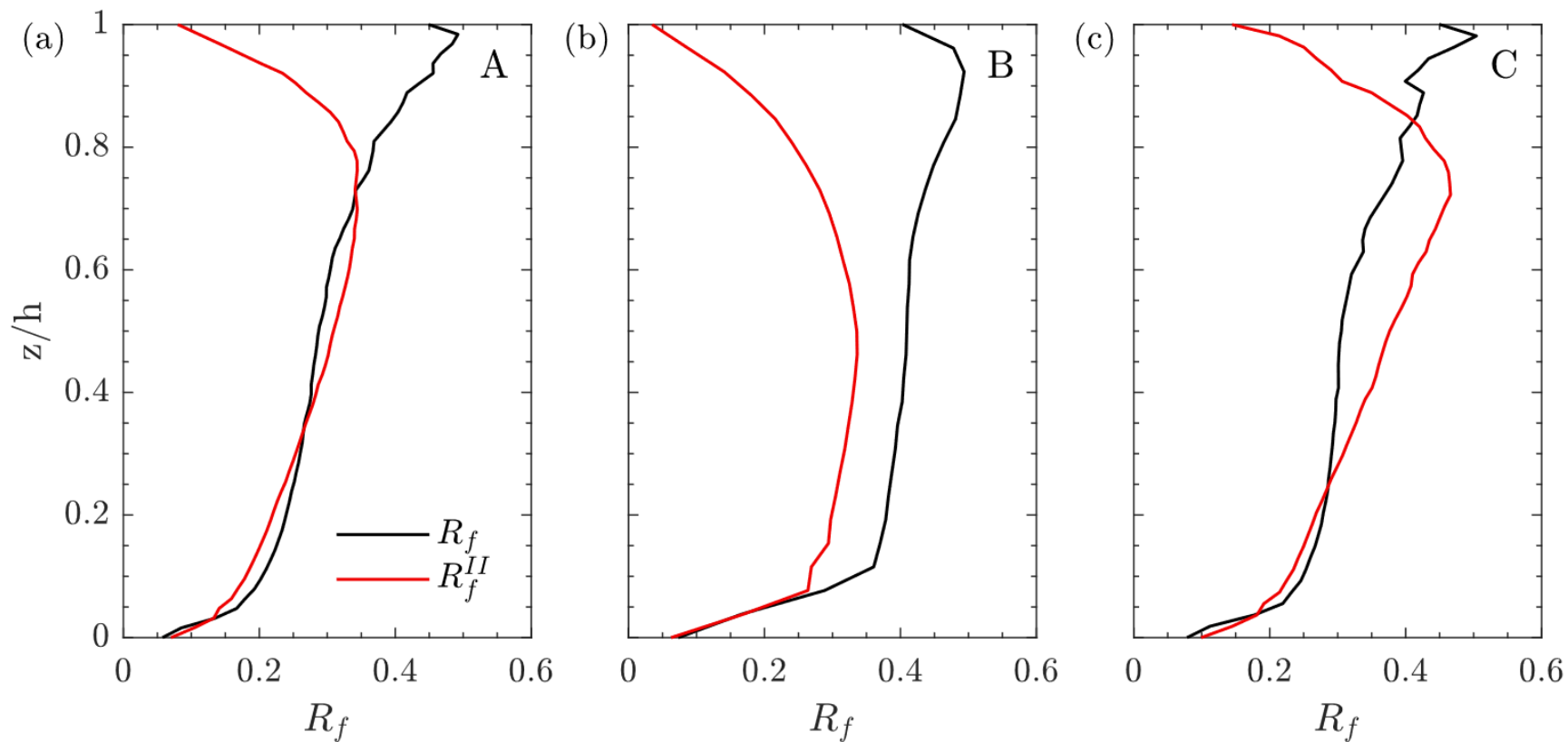
D – destruction of TKE by buoyancy and dissipation



Mixing efficiency

$$R_f = \frac{B}{P}$$

$$R_f^{II} = \frac{B}{m} = \frac{B}{B + \varepsilon}$$



Conclusions

- Successful large-eddy simulations of very stable conditions ($h/L > 6$ and $Ri > 0.25$).
- LES is very computationally challenging because of very small time-steps and long-time integrations.
- Wind turning significant from the surface.
- Turbulent flows met the criteria of strongly stratified regime without viscous effects ($Re_b \gg 1$ and $Fr_h \ll 1$), except near the SBL top where the flow is viscously-dominated.
- Flow is very anisotropic.

Mixing efficiency

