Eddy-Diffusivity Mass Flux Addates and Edd ARPEGE

Secure loup CNRM/GAME Météo-France and CNRS

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Working days on EFB closure 18-22 march 2013





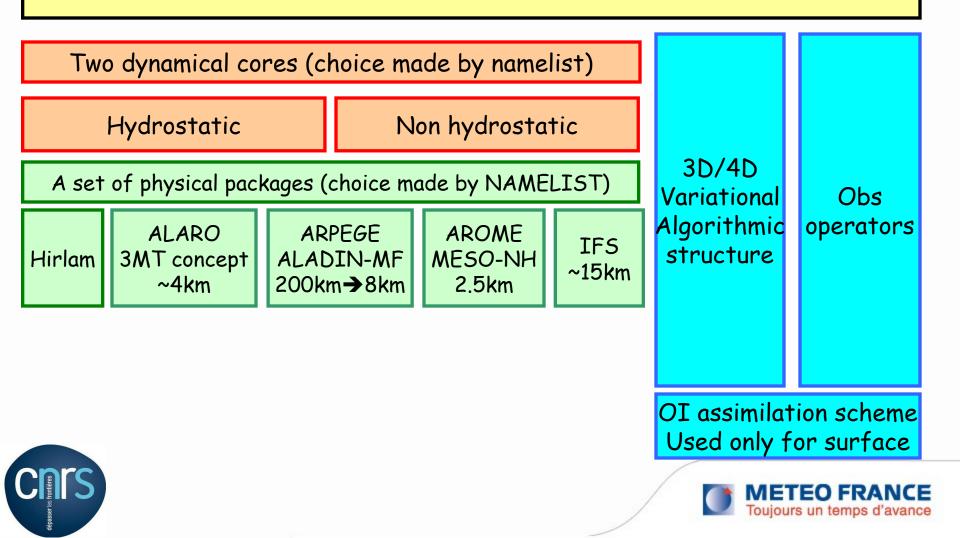
- ➤ The ARPEGE/IFS/AROME/HARMONY... world
- Current operational situation (ARPEGE and AROME)
- EDMF concept
- Evolution strategy (seamless approach, convergence with AROME)
- Stability problem and implicit solution
- > References





ARPEGE/ALADIN/AROME/IFS/HARMONIE A unified sofware

GLOBAL (variable mesh or not) or LAM (choice made by NAMELIST)



> ARPEGE is a global spectral model with a variable mesh

> T798 C=2.4 (Δt = 514s) \rightarrow 10 km over France and around 60 km at the antipode, few hundred kilometers east New-Zealand

➢ 70 vertical levels → Close to ECMWF vertical resolution in the troposphere

> 4DVAR multi-incremental data assimilation, with two outer loops T107 C=1 (Δt = 1800s) and T323 C=1 (Δt = 1350s) using a 6 hours window

> ALADIN-MF is an hydrostatic LAM with the same physics it runs over Indien Ocean, West Indies, French Polynesia, New-Caledonia and some secret parts of the world (army queries !)

> 3DVAR data assimilation



> Presently 8km, 70 levels, $\Delta t = 480s$



- > AROME is a non-hydrostatic LAM
- Physical parametrizations come from Méso-Nh
- >It runs over France (coupling model is ARPEGE)
- > 3DVAR data assimilation
- Presently 2.5km, 60 levels (more levels than ARPEGE in the PBL)
- ≻ ∆t = 60s





Operationnal «NWP» Boundary layer physics at Météo-France

 $K = cL_{BL89}\sqrt{TKE}$

lacarrère (1989)

Where I_{up} and I_{down} are computed using

dry buoyancy following Bougeault and

All NWP models (AROME, ARPEGE and ALADIN-MF) use « EDMF » concept (Hourdin et al 2002, Soares et al 2004, Siebesma et al 2007)

$$\overline{w'\phi'} = -K \frac{\partial \phi}{\partial z} + \frac{M_u}{\rho} \left(\phi_u - \overline{\phi} \right) \quad \text{with}$$
and
$$L_{BL89} = \begin{bmatrix} \frac{(l_{up})^{-\frac{2}{3}} + (l_{down})^{-\frac{2}{3}}}{2} \end{bmatrix}^{-\frac{3}{2}} \quad \text{When}$$

$$dry \text{ b}$$

$$lacar$$

ARPEGE and ALADIN-MF

Prognostic turbulent kinetic energy scheme « CBR » (Cuxart et al 2000)

Shallow convection mass flux scheme « KFB » (Bechtold et al 2001) Equations should be the same AROME

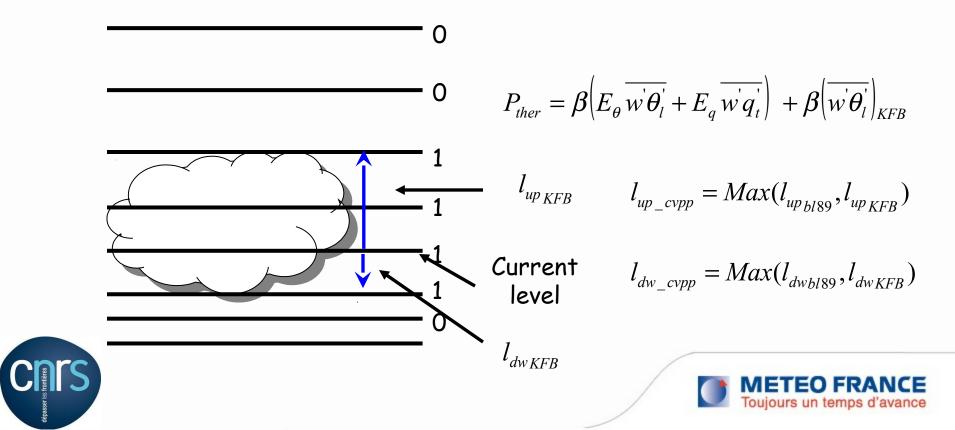
Prognostic turbulent kinetic energy scheme « CBR » (Cuxart et al 2000)

 Shallow convection and dry thermal mass flux scheme
 « EDKF » (Pergaud et al 2009)

dána

> With KFB, during our first evaluation tests in ARPEGE, we found too much low level clouds and too much wind in the PBL in the tropical area

> A thermal production term is then computed by KFB and Bougeault Lacarrère (1989) mixing lengths are increased in the shallow clouds



It was found a large beneficial impact on wind in the tropics (205 \rightarrow 20N)

- 71 750 🗲 🚥 Zonal mean over the tropical area of the Kinetic energy (J/kg)-000 (hPa) with (red) and without (black) the thermal production term 850 → 50prag coming from shallow convection and the modification of the mixing -800length inside the cloud. 950 → 💷



Crantos

- > No dry thermal in KFB
- > No mixing of wind in KFB
- Convergence strategy between NWP models physics (seamless approach)
- > Global model is a great testbed for parametrizations
- > But, global models are very sensitive clockworks
- > KFB is numerically stable at large time step \rightarrow T107 $\Delta t = 1800s$
- > With EDKF we uncountered numerical stability problems

> The solution was a common implicit solver for Eddy-Difusivity and Mass Flux part





Implicit treatment of the Mass Flux equation (1)

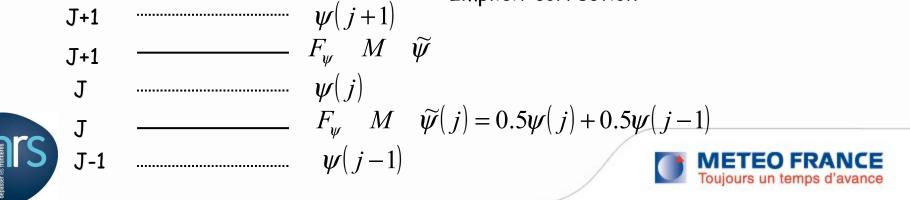
$$\begin{cases} F_{\psi} = \rho \overline{w'\psi'} = M(\psi_u - \overline{\psi}) \\ \left(\frac{\partial \psi}{\partial t}\right)_{MF} = \frac{1}{\rho} \frac{\partial}{\partial z} F_{\psi} \end{cases}$$

Second equation is solved implicitly $F_{\psi} = (1 - z_i)F_{\psi}^- + z_iF_{\psi}^+$

$$F_{\psi}^{+} = F_{\psi}^{-} + \delta F_{\psi} = F_{\psi}^{-} + \frac{\partial F_{\psi}}{\partial \psi} \delta \psi = F_{\psi}^{-} - M \left(\widetilde{\psi}^{+} - \widetilde{\psi}^{-} \right)$$

Then :
$$\left(\frac{\partial \psi}{\partial t}\right)_{MF} = \frac{1}{\rho} \frac{\partial}{\partial z} \left(F_{\psi}^{-} - \underbrace{z_{i} M\left(\widetilde{\psi}^{+} - \widetilde{\psi}^{-}\right)}_{MF}\right)$$

Implicit correction



We obtain :
$$\begin{split} \psi^{+}(j) - \psi^{-}(j) &= \frac{\Delta t}{\rho \Delta z} \Big[F_{\psi}^{-}(j+1) - F_{\psi}^{-}(j) \\ &- z_{i} M(j+1) \Big| 0.5 \, \psi^{+}(j+1) + 0.5 \, \psi^{+}(j) - 0.5 \, \psi^{-}(j+1) - 0.5 \, \psi^{-}(j) \Big| \\ &+ z_{i} M(j) \Big| 0.5 \, \psi^{+}(j) + 0.5 \, \psi^{+}(j-1) - 0.5 \, \psi^{-}(j) - 0.5 \, \psi^{-}(j-1) \Big| \Big] \end{split}$$

Grouping '+' terms in the left hand side of the equation we obtain the following tridiagonal system :

$$\psi^{+}(j+1) \left[0.5 \frac{\Delta t}{\rho \Delta z} z_{i} M(j+1) \right]$$

$$+ \psi^{+}(j) \left[1 + 0.5 \frac{\Delta t}{\rho \Delta z} z_{i} M(j+1) - 0.5 \frac{\Delta t}{\rho \Delta z} z_{i} M(j) \right]$$

$$- \psi^{+}(j-1) \left[0.5 \frac{\Delta t}{\rho \Delta z} z_{i} M(j) \right] = \psi^{-}(j) + \frac{\Delta t}{\rho \Delta z} \left(F_{\Psi}^{-}(j+1) - F_{\Psi}^{-}(j) \right)$$

$$+ 0.5 \frac{\Delta t}{\rho \Delta z} z_{i} M(j+1) \left(\psi^{-}(j+1) + \psi^{-}(j) \right)$$

$$- 0.5 \frac{\Delta t}{\rho \Delta z} z_{i} M(j) \left(\psi^{-}(j) + \psi^{-}(j-1) \right)$$

$$\longrightarrow$$
The second second

Implicit treatment of the Eddy-Difusivity equation

Eddy Difusivity equation,

$$\left(\frac{\partial\psi}{\partial t}\right)_{eddy} = -\frac{1}{\rho}\frac{\partial}{\partial z}\left(k\frac{\partial\psi}{\partial z}\right)$$

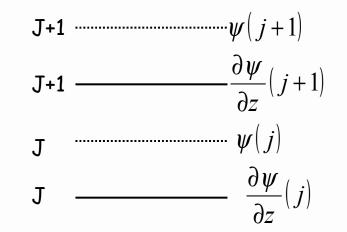
is discretized as follows :

$$\psi^{+}(j) - \psi^{-}(j) = -\frac{\Delta t}{\rho \Delta z(j)} \left[\frac{k(j+1)}{\Delta z(j+1)} \left(\psi^{+}(j+1) - \psi^{+}(j) \right) - \frac{k(j)}{\Delta z(j)} \left(\psi^{+}(j) - \psi^{+}(j-1) \right) \right]$$

This yields to the simple tridiagonal system :

$$\psi^{+}(j+1) \left[\frac{\Delta t}{\rho \Delta z(j)} \frac{k(j+1)}{\Delta z(j+1)} \right]$$

+ $\psi^{+}(j) \left[1 - \frac{\Delta t}{\rho \Delta z(j)} \left(\frac{k(j+1)}{\Delta z(j+1)} + \frac{k(j)}{\Delta z(j)} \right) \right]$
+ $\psi^{+}(j-1) \left[\frac{\Delta t}{\rho \Delta z(j)} \frac{k(j)}{\Delta z(j)} \right] = \psi^{-}(j)$





Common implicite resolution of the EDMF equation

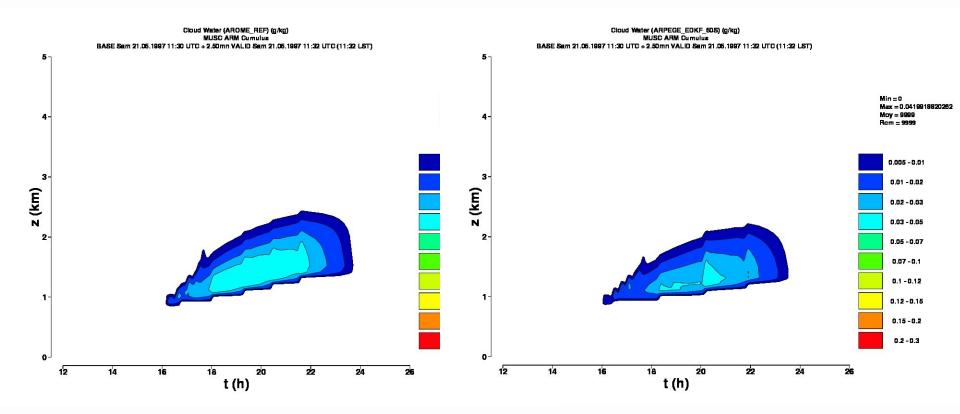
Discretization of the full EDMF equation :

$$\left(\frac{\partial\psi}{\partial t}\right)_{edmf} = \frac{1}{\rho}\frac{\partial}{\partial z}\left(-k\frac{\partial\psi}{\partial z} + M(\psi_u - \overline{\psi})\right)$$

Yields to the following tridiagonal system :

$$\begin{split} \psi^{+}(j+1) \Bigg[\frac{\Delta t}{\rho \Delta z(j)} \bigg(\frac{k(j+1)}{\Delta z(j+1)} + 0.5M(j+1) \bigg) \Bigg] \\ + \psi^{+}(j) \Bigg[1 - \frac{\Delta t}{\rho \Delta z(j)} \bigg(\frac{k(j+1)}{\Delta z(j+1)} + \frac{k(j)}{\Delta z(j)} + 0.5M(j+1) - 0.5M(j) \bigg) \Bigg] \\ + \psi^{+}(j-1) \Bigg[\frac{\Delta t}{\rho \Delta z(j)} \bigg(\frac{k(j)}{\Delta z(j)} + 0.5M(j) \bigg) \Bigg] = \psi^{-}(j) + \frac{\Delta t}{\rho \Delta z(j)} \Big(F_{\psi}^{-}(j+1) - F_{\psi}^{-}(j) \Big) \\ + 0.5 \frac{\Delta t}{\rho \Delta z(j)} M(j+1) \Big(\psi^{-}(j+1) + \psi^{-}(j) \Big) \\ - 0.5 \frac{\Delta t}{\rho \Delta z(j)} M(j) \Big(\psi^{-}(j) + \psi^{-}(j-1) \Big) \Big] \end{split}$$

Test in 1D model using the Arm Cumulus case (1)



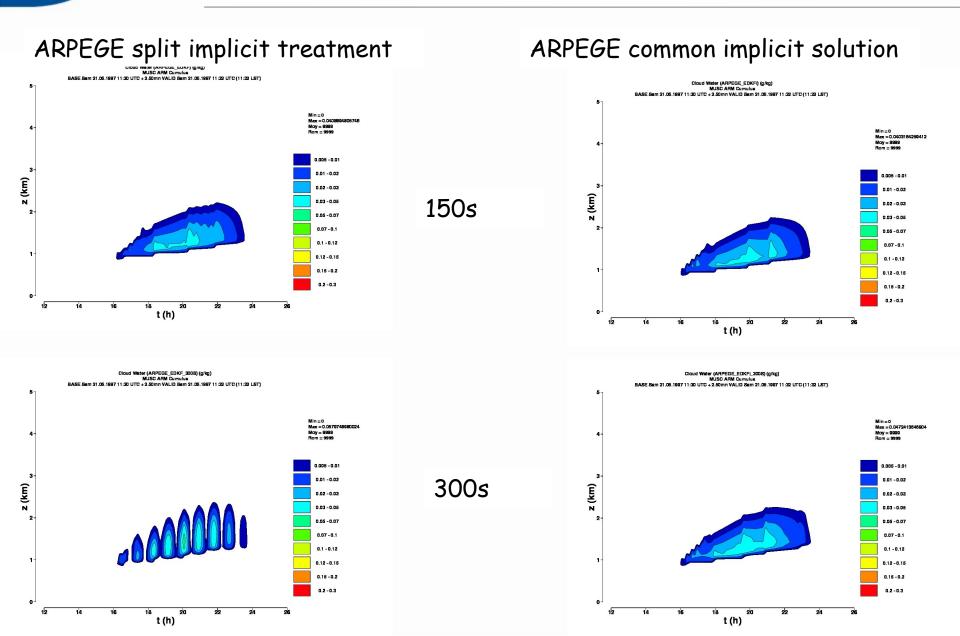
Cloud liquid water AROME 60s

Cloud liquid water ARPEGE 60s





Test in 1D model using the Arm Cumulus case (2)





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